

# Homework #6 Solutions

1. (a) Threshold voltage of a symmetric, fully-depleted dual-gate MOSFET

By symmetry,  $\epsilon = 0$  in center of thin Si region (at  $x = x_s/2$ ).

Both channels are equivalent (turn on at same time) so focus on left ( $x=0$ )  $\leftarrow$  net doping

from center of thin body  $\rightarrow$

$$\psi_s = \frac{(x_s/2)^2 q N_A}{2\epsilon_s \epsilon_0}$$

$$V_{ox} = -\frac{Q_s/2}{C_{ox}} = +\frac{q N_A x_s/2}{C_{ox}}$$

(half of  $Q_s$  supported by each gate) at  $\epsilon_0 \epsilon_{Si}$ ,  $\psi_I \approx 0.5V$

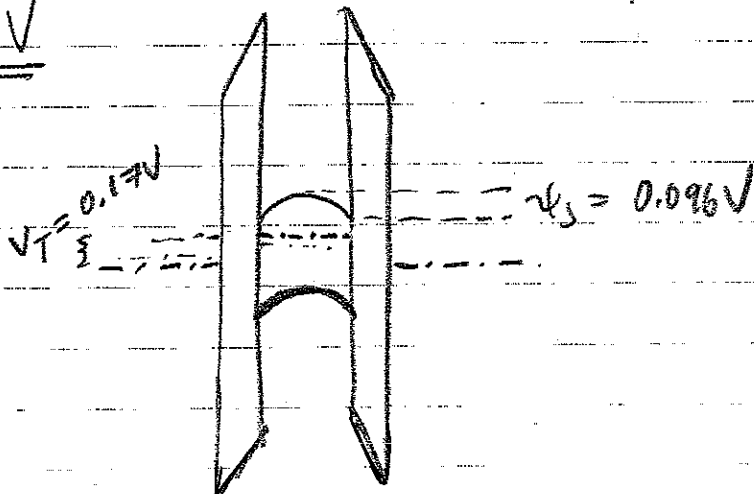
$$V_T = \Phi_{MSi} + V_{ox}$$

$$\Phi_{MSi} = \Phi_M - \Phi_{SE} = \Phi_M - \left( x_s + \frac{\epsilon_s}{2q} - \psi_I \right)$$

$$V_T = \Phi_M - 4.11V + \frac{q N_A x_s/2}{2K_{ox}\epsilon_0}$$

$$(b) V_T = 4.05 - 4.11 + \frac{q (2 \times 10^{19} \text{ cm}^{-3}) (5 \times 10^{-7} \text{ cm}) (10^{-7} \text{ cm})}{2 (3.9) (8.854 \times 10^{-14} \text{ F/cm})} = \underline{\underline{0.17V}}$$

$$\psi_s = 0.097$$



2(a) From Eq (18)

$$\frac{d\langle p \rangle}{dE} = -\frac{1}{m^*} \langle p \rangle \cdot \nabla_r \langle \hat{p} \rangle - qE - \frac{1}{\hbar} \nabla_r (\hbar k T_e) - \frac{\langle \hat{p} \rangle}{\langle \tau_m \rangle}$$

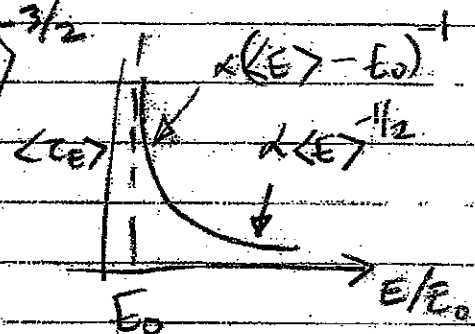
steady-state
Spatially uniform

$$\langle \hat{p} \rangle = -qE \langle \tau_m \rangle = -qE \left[ A \langle E \rangle^{3/2} \right]$$

$$|v_d| = \frac{|\langle p \rangle|}{m^*} = \frac{v_{sat} E}{E + E_c} = \frac{qE}{m^*} A \langle E \rangle^{3/2}$$

$$\langle E \rangle^{-3/2} = \frac{v_{sat} m^*}{qA(E + E_c)}$$

$$\langle E \rangle = \left( \frac{qA(E + E_c)}{v_{sat} m^*} \right)^{2/3}$$



(b) Use Eq. (19) to get  $0 = -\frac{1}{m^*} qE \langle p \rangle - \frac{\langle E \rangle - E_0}{\langle \tau_m \rangle}$

$$+ qE \left( \frac{qE}{m^*} A \langle E \rangle^{3/2} \right) = \frac{\langle E \rangle - E_0}{\langle \tau_m \rangle}$$

$$\langle \tau_m \rangle = \frac{(\langle E \rangle - E_0) m^* \langle E \rangle^{3/2}}{q^2 E^2 A} = \frac{(\langle E \rangle - E_0) m^* \langle E \rangle^{3/2}}{q^2 E^2 A}$$

From above,

$$E = \frac{v_{sat} m^* \langle E \rangle^{3/2}}{qA} - E_c$$

$$\text{For } E=0, \bar{E} = E_0, \text{ so } E_0 = \left( \frac{qA E_c}{v_{sat} m^*} \right)^{2/3}; E_c = \frac{v_{sat} m^*}{qA} E_0^{3/2}$$

$$\langle \tau_m \rangle = \frac{(\langle E \rangle - E_0) m^* \langle E \rangle^{3/2}}{\left( \frac{v_{sat} m^*}{qA} \right)^2 \left( \langle E \rangle^{3/2} - E_0^{3/2} \right)^2} \quad \text{See sketch above}$$

$$\beta = \frac{4 (3kT m^*) \epsilon_s kT}{\hbar^2 q^2 N_I} = \frac{12 (kT/q)^2 m^* \epsilon_s}{\hbar^2 N_I}$$

$$= \frac{12 (0.026 \text{ V})^2 (0.33 \times 9.11 \times 10^{-31} \text{ kg}) (8.854 \times 10^{-14} \frac{\text{F}}{\text{cm}}) 12}{(1.05 \times 10^{-34} \text{ J}\cdot\text{s})^2 (10^{16} \text{ cm}^{-3})} \quad (1 \text{ eV} = 1 \text{ V})$$

$$= 2.4 \times 10^7 \frac{\text{V}^2 \text{ kg} \frac{\text{C}}{\text{V}\cdot\text{cm}}}{\text{J}^2 \text{ s}^2 \text{ cm}^{-3}} = 2.4 \times 10^7 \frac{\text{kg}}{\text{J}^2 \text{ s}^2 \text{ cm}^{-3}} = 2400$$

$$\beta = 2400 \left( \frac{T/300\text{K}}{(N_I/10^{16} \text{ cm}^{-3})} \right)^2, \quad \mu = \frac{2.2 \times 10^4 (T/300\text{K})^{3/2}}{(N_I/10^{16} \text{ cm}^{-3}) \left[ \ln|1+\beta| - \frac{\beta}{1+\beta} \right]}$$

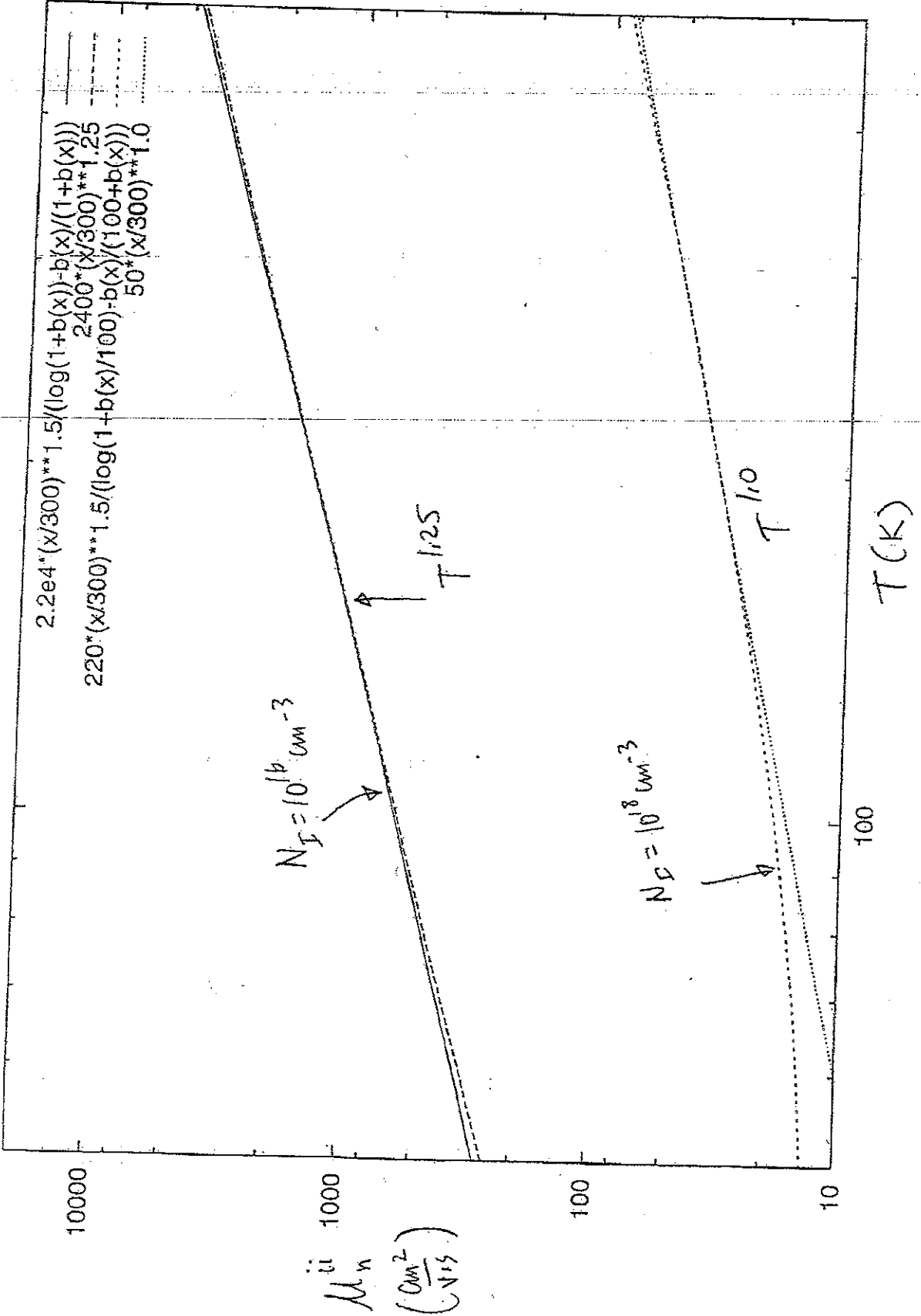
Looking at plot on next page; near 300K,  $\mu \propto T^{1.25}$  for  $N_I = 10^{16}$   
 $\mu \propto T$  for  $N_I = 10^{18}$

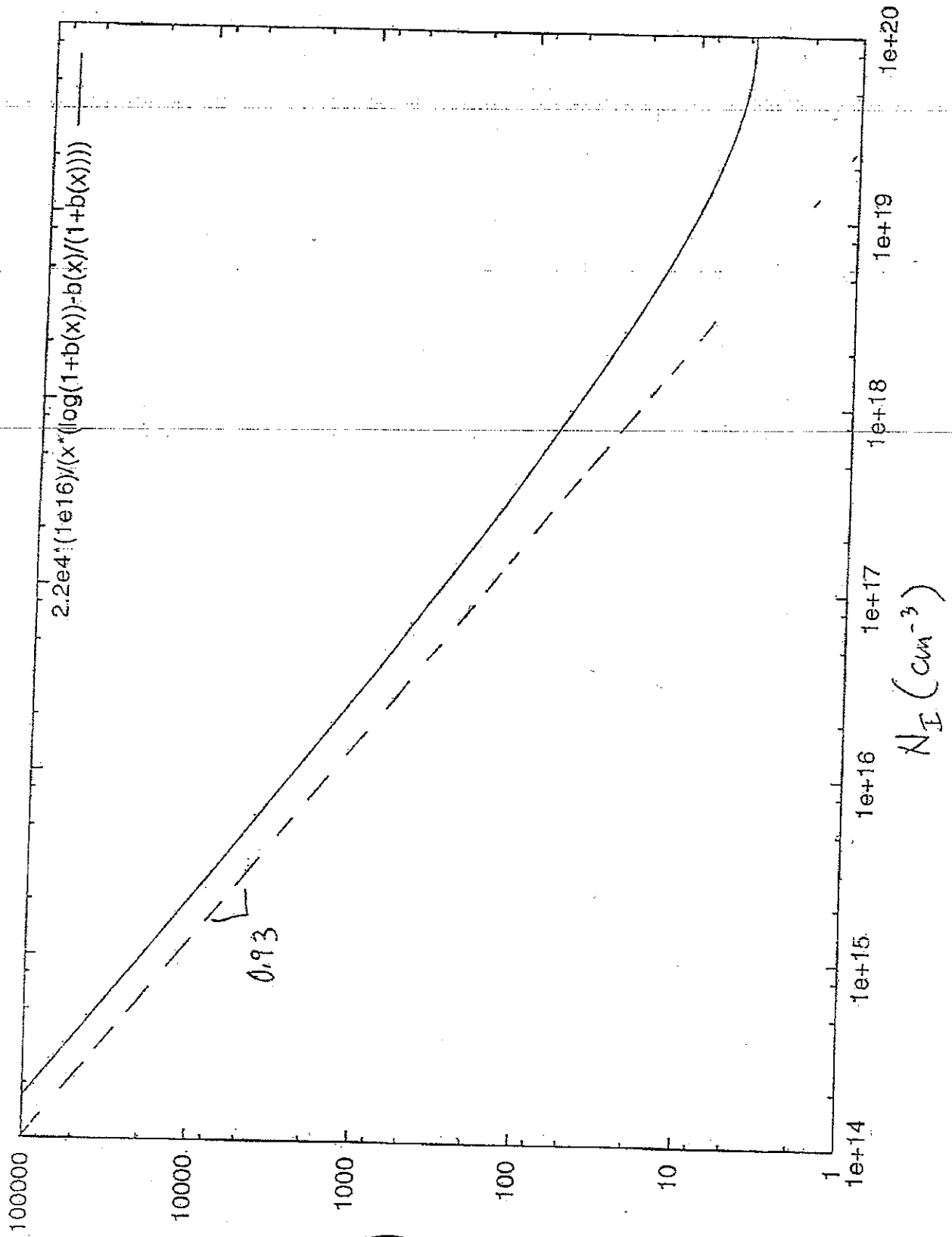
As expected, mobility is limited by ionized impurity scattering increases with T due to high speed of electrons. The power dependence varies with T and doping level. The dependence is slightly weaker than the simple  $T^{3/2}$  dependence predicted by a simple model (leading term).

(b) Plot is shown on following page based on  $\beta(T, N_I), \mu(T, N_I)$  above

As expected, mobility drops with increased doping. The power dependence is slightly weaker than linear at low  $N_I$ , and saturates at high  $N_I$ .

Note that at low  $N_I$  or high T,  $\mu$  will be limited by phonon rather than ionized impurity scattering.





$$\alpha_n \left( \frac{\text{cm}^2}{\text{V} \cdot \text{s}} \right)$$

3

$$\mu = \frac{q Z_m}{m^*}$$

$$\frac{1}{Z_m} = \frac{\pi N_I Z^2 q^4 N(E_k)}{\hbar \epsilon_s^2 k^4} \left[ \ln \left| 1 + \frac{4k^2}{\hbar_D^2} \right| - \frac{4k^2/\hbar_D^2}{1 + 4k^2/\hbar_D^2} \right]$$

$$N(E_k) = \frac{4\pi}{\hbar^3} (2m^*)^{3/2} E_k, \text{ where } m^* = (m_1 m_2)^{1/2} \text{ and } E_k = E - E_c = 0.33 m_0$$

$$\frac{\hbar_D^2}{2} = \frac{q^2 n_0}{\epsilon_s k T}, \quad k^2 = \frac{2 E_k m^*}{\hbar^2}, \quad Z = 1 \text{ (single ionized)}$$

a) Assume  $n_0 = N_I$ , and that on average  $E_k = \frac{3}{2} kT$

Substituting above, 
$$\mu = \frac{q \hbar \epsilon_s^2 k^4 \hbar^3}{m^* \pi N_I Z^2 q^4 4\pi (2m^*)^{3/2} E_k^{1/2} \cdot \left[ \ln \left| 1 + \frac{\beta}{1+\beta} \right| - \frac{\beta}{1+\beta} \right]}$$

where 
$$\beta = \frac{4k^2}{\hbar_D^2} = \frac{4 E_k m^*}{\hbar^2 \epsilon_s^2 (kT)^2}$$

$$\beta = \frac{4 (2 E_k m^*) \epsilon_s^2 (kT)^2}{\hbar^2 q^2 N_I} \quad \mu = \frac{\epsilon_s^2 4 E_k^{3/2}}{m^{*1/2} N_I q^3 (2\pi)^{3/2} E_k^{1/2} \cdot \left[ \dots \right]}$$

$$\mu = \frac{\epsilon_s^2 \pi (3kT)^{3/2} (2\pi)}{m^* N_I q^3 2^3 \left[ \dots \right]} = \frac{\epsilon_s^2 3^{3/2} (kT)^{3/2} \pi}{m^* N_I q^3 \left[ \dots \right]} \quad \left( \frac{1}{2} \right)^{1/2} = \left( \frac{q m^*}{\epsilon_s^2} \right)^{1/2}$$

At  $T = 300K$ ,  $N_I = 10^{16} \text{ cm}^{-3}$

$$\mu = \frac{12 (8.854 \times 10^{-14} \text{ F/cm})^2 (3(0.026 \text{ V}))^{3/2} \pi}{(0.33 \times 9.11 \times 10^{-31} \text{ kg})^{1/2} (10^{16} \text{ cm}^{-3}) (1.6 \times 10^{-19} \text{ C})^3} = \frac{220 \frac{\text{C}^2}{\text{V}^2 \text{cm}^2} \text{V}^{3/2}}{\text{kg}^{1/2} \text{cm}^{-3} \text{C}^{3/2}} = 2.2 \times 10^4 \frac{\text{cm}^2}{\text{V} \cdot \text{s}}$$